

REMARK ON THE UNCONDITIONAL STABILITY OF A FINITE ELEMENT SCHEME FOR NONLINEAR PLATE PROBLEMS

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1. Introduction

The nonlinear vibration of a thin elastic plate can be represented by a system of two nonlinear partial differential equations. For solving this system by the finite element method, one of the present authors derived a stability criterion with respect to the time increment Δt under a certain condition. The purpose of this note is to show that this condition is unnecessary to ensure the stability.

We consider the equations in the cylindrical domain $Q = (0, T) \times \Omega$, being Ω the shape of the plate. Let f be the Airy's stress function and w the deflection of the plate. Then the system of equations is

$$(1.1) \quad \begin{cases} \Delta^2 f = -[w, w] \\ w_{,tt} - \Delta w_{,tt} + \Delta^2 w = [f + f_0, w] + p, \end{cases}$$

where Δ^2 and $[f, w]$ denote the biharmonic operator and $f_{,xx} w_{,yy} + f_{,yy} w_{,xx} - 2f_{,xy} w_{,xy}$, respectively. The functions p and f_0 correspond to given lateral load and a stress function derived from given plain-stress problem, respectively. We assume that the boundary $\partial\Omega$, p and f_0 are sufficiently smooth. The system (1.1) is solved under the boundary condition $f = w = \frac{df}{dn} = \frac{dw}{dn} = 0$ on $\partial\Omega$ and the initial condition $w|_{t=0} = w_0, \frac{\partial w}{\partial t}|_{t=0} = w_1$, where n is the outward normal to $\partial\Omega$.

2. Approximating scheme.

We use (u, w) and $\|u\|$ to denote the inner product and the norm in $L_2(\Omega)$. The space H is the completion of the set of all C^∞ -functions with support in Ω under the norm $\|w\|_H = \sqrt{(\Delta w, \Delta w)}$. By \hat{H} we denote the finite element subspace of H spanned by a finite element basis $\{\phi_i\}$ ($i=1, \dots, K$).

The approximating scheme presented in [1] is as follows. The interval $[0, T]$ is divided into equal pieces of length $\Delta t = T/N$ by the points $t = n\Delta t$ ($n=0, 1, \dots, N$).

Then the approximate solution (f^n, w^n) at the time level $n\Delta t$ is determined by the following system of equations in \hat{H} .

$$(2.1) \quad (\Delta f^n, \Delta \phi) = -\frac{1}{2} ([w^{n+1} + w^{n-1}, w^n], \phi) \quad \text{for all } \phi \in \hat{H},$$

$$(2.2) \quad \begin{aligned} (D_t D_i w^n, \phi) - (D_t D_{\bar{i}} \Delta w^n, \phi) + \frac{1}{2} (\Delta(w^{n+1} + w^{n-1}), \Delta \phi) \\ = ([f^n + f_0^n, w^n], \phi) + (p^n, \phi) \quad \text{for all } \phi \in \hat{H}, \end{aligned}$$

where D_t and $D_{\bar{i}}$ are the forward and backward difference operators, respectively. The initial conditions are approximated as

$$(2.3) \quad w^0 = \dot{w}_0, \quad w^1 = \dot{w}_0 + \dot{w}_1 \Delta t,$$

where \dot{w} denotes the interpolate of w in \hat{H} . This scheme is well defined. Precisely, the coefficient matrix to determine w^{n+1} is symmetric and positive definite through all time steps (see [1]).

To analyze the approximating scheme it is convenient to represent it by a system of operator equations in H .

We define the following operators C and B by means of the Riesz representation theorem for bounded linear functional on H :

$$\begin{aligned} ([u, v], \phi) &= (C(u, v), \phi)_H \quad \text{for all } \phi \in H, \\ (u, \phi) &= (Bu, \phi)_H \quad \text{for all } \phi \in H. \end{aligned}$$

Let P be the projection on H onto \hat{H} . Then our scheme is represented on H as follows.

$$(2.4) \quad \begin{aligned} f^n &= -\frac{1}{2} PC(w^{n+1} + w^{n-1}, w^n) \\ PB(D_t D_i w^n - D_t D_{\bar{i}} \Delta w^n) &+ \frac{1}{2} (w^{n+1} + w^{n-1}) \\ &= PC(f^n + f_0^n, w^n) + PBp^n. \end{aligned}$$

As well known, the form $([u, v], w)$ is symmetric, that is,

$$(2.5) \quad ([u, v], w) = [v, u], w) = ([w, u], v).$$

3. Stability in energy.

For proving the unconditional stability of our scheme, we provide the following lemma.

LEMMA 1. Let $\{x_n\} (n=0, 1, \dots, N)$ be a sequence of nonnegative numbers satisfying the following inequality.

$$D_t x_n \leq C_1 x_n^{\frac{1}{2}} + C_2 x_{n-1}^{\frac{1}{2}},$$

where the constants C_1, C_2 are positive, and x_0 is given. Then, holds the following inequality for all n .

$$x_n^{\frac{1}{2}} \leq x_0^{\frac{1}{2}} + (C_1 + C_2)T.$$

PROOF. Consider the positive solution of the equation,

$$D_t x'_n = C_1 x'_n{}^{\frac{1}{2}} + C_2 x'_{n-1}{}^{\frac{1}{2}}, \quad x'_0 = x_0.$$

This equation has a unique nonnegative solution $\{x'_n\}$ satisfying

$$(3. 1) \quad x_n \leq x'_n \text{ and } x_0 = x'_0 \leq x'_1 \leq x'_2 \leq \dots \leq x'_N.$$

Therefore, the following inequality holds for the sequence $\{x'_n\}$.

$$D_t x'_n \leq (C_1 + C_2) x'_n{}^{\frac{1}{2}}.$$

We next consider the equation obtained by equating the both sides of this inequality:

$$D_t x''_n = C x''_n{}^{\frac{1}{2}} \quad \text{where} \quad C = C_1 + C_2, \quad x''_0 = x'_0.$$

This equation has a unique nonnegative solution $\{x''_n\}$ satisfying,

$$(3. 2) \quad x'_n \leq x''_n \text{ and } x'_0 = x''_0 \leq x''_1 \leq x''_2 \leq \dots \leq x''_N.$$

It is easy to see that $x''_n{}^{\frac{1}{2}} \leq x''_{n-1}{}^{\frac{1}{2}} + CAt$. Hence we have

$$(3. 3) \quad x''_n{}^{\frac{1}{2}} \leq x''_0{}^{\frac{1}{2}} + CT.$$

The lemma follows from (3.1), (3.2) and (3.3).

Our conclusion is then as follows.

THEOREM 1. Let E_n be defined by

$$E_n = \|D_t w^n\|^2 + |D_t w^n|_1^2 + \frac{1}{2} \|Aw^{n+1}\|^2 + \frac{1}{2} \|Aw^n\|^2 + \frac{1}{4} \|PC(w^n, w^{n+1})\|_H^2 \\ + \frac{1}{4} \|PC(w^n, w^{n+1}) - 2f_0^n\|_H^2 - \|f_0^n\|_H^2 + M,$$

where $M = \text{Max}_t \|f_0\|_H^2$, and $|w|_1^2 = \|w_x\|^2 + \|w_y\|^2$. Then, holds the following energy inequality.

$$E_n^{\frac{1}{2}} \leq E_0^{\frac{1}{2}} + CT \quad (n=1, 2, \dots, N),$$

where $C = 2\text{Max}_t (\|D_t f_0\|_H + \|\dot{p}\|)$.

PROOF. Replacing ϕ by $(w^{n+1} - w^{n-1})/\Delta t$ in the both sides of the equation (2. 2), we have the next equality.

$$(3. 4) \quad (D_t D_t w^n, w^{n+1} - w^{n-1})/\Delta t - (D_t D_t \Delta w^n, w^{n+1} - w^{n-1})/\Delta t \\ + \frac{1}{2} (\Delta(w^{n+1} + w^{n-1}), \Delta(w^{n+1} - w^{n-1}))/\Delta t \\ = ([f^n + f_0^n, w^n], w^{n+1} - w^{n-1})/\Delta t + (\dot{p}^n, w^{n+1} - w^{n-1})/\Delta t.$$

Each term of the left side is written as follows.

$$(D_t D_t w^n, w^{n+1} - w^{n-1})/\Delta t = D_t \|D_t w^n\|^2, \\ - (D_t D_t \Delta w^n, w^{n+1} - w^{n-1})/\Delta t = D_t |D_t w^n|_1^2, \\ \frac{1}{2} (\Delta(w^{n+1} + w^{n-1}), \Delta(w^{n+1} - w^{n-1}))/\Delta t = \frac{1}{2} D_t (\|Aw^{n+1}\|^2 + \|Aw^n\|^2).$$

On the other hand, by using the first equality of (2. 4) and the relation (2. 5), we have

$$([f^n + f_0^n, w^n], w^{n+1} - w^{n-1})/\Delta t \\ = -\frac{1}{2} D_t \|PC(w^n, w^{n+1})\|_H^2 + ([f_0^n, w^n], w^{n+1} - w^{n-1})/\Delta t \\ = -\frac{1}{2} D_t \|PC(w^n, w^{n+1})\|_H^2 + D_t (PC(w^n, w^{n+1}), f_0^n)_H \\ - (PC(w^n, w^{n-1}), D_t f_0^n)_H \\ = -D_t \left[\frac{1}{4} \|PC(w^n, w^{n+1})\|_H^2 + \frac{1}{4} \|PC(w^n, w^{n+1}) - 2f_0^n\|_H^2 - \|f_0^n\|_H^2 \right] \\ - (PC(w^n, w^{n-1}), D_t f_0^n)_H.$$

Since $(\dot{p}^n, w^{n+1} - w^{n-1})/\Delta t = (\dot{p}^n, D_t w^n + D_t w^{n-1})$, the equality (3. 4) is written as follows.

$$\begin{aligned}
D_i & [\|D_i w^n\|^2 + |D_i w^n|_1^2 + \frac{1}{2} \|w^{n+1}\|^2 + \frac{1}{2} \|Aw^n\|^2 + \frac{1}{4} \|PC(w^n, w^{n+1})\|_H^2 \\
& + \frac{1}{4} \|PC(w^n, w^{n+1}) - 2f_0^n\|_H^2 - \|f_0^n\|_H^2] \\
& = - (PC(w^n, w^{n+1}), D_i f_0^n)_H + (p^n, D_i w^n + D_i w^{n-1}).
\end{aligned}$$

Since M is constant, we have

$$\begin{aligned}
D_i E_n & \leq \|PC(w^n, w^{n-1})\|_H \|D_i f_0^n\|_H + \|p^n\| (\|D_i w^n\| + \|D_i w^{n-1}\|) \\
& \leq (2\|D_i f_0^n\|_H + \|p^n\|) E_{n-1}^{\frac{1}{2}} + \|p^n\| E_n^{\frac{1}{2}}.
\end{aligned}$$

Hence, by the lemma, we have the following estimation.

$$E_n^{\frac{1}{2}} \leq E_0^{\frac{1}{2}} + CT,$$

where $C = 2 \underset{t}{\text{Max}} (\|D_i f_0\|_H + \|p\|)$. The theorem is thus proved.

REMARK: In [1], the stability is proved under the condition

$$\Delta t < 2 / (3\sqrt{3}) \underset{|\alpha|=2}{\text{Max}} |D^\alpha f_0^n|.$$

Reference

- [1] T. Miyoshi: An unconditionally stable implicit finite element scheme for solving nonlinear dynamical problems of elastic plates, FUNCTIONAL ANALYSIS AND NUMERICAL ANALYSIS. Japan-France Seminar, Tokyo and Kyoto, 1976; H. Fujita (Ed.), Japan Society for the Promotion of Science, 1978.

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